

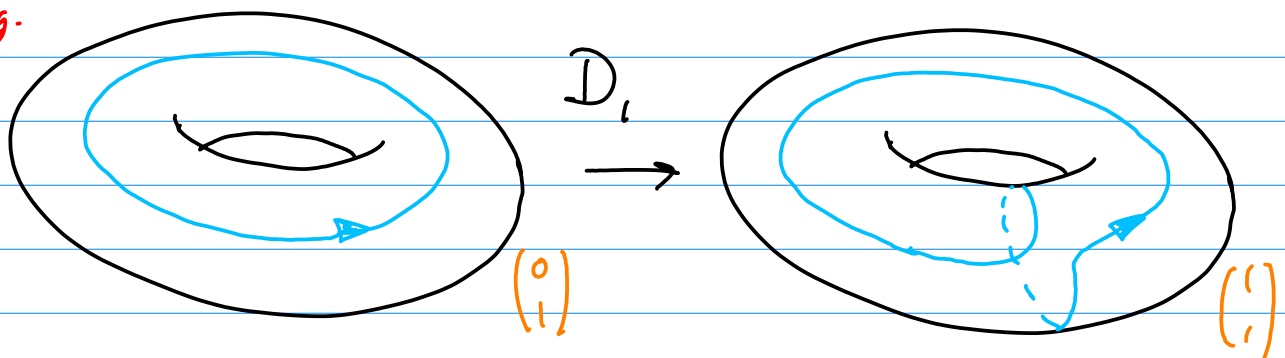
# Braids Lectures 13-14: Mapping Class Group of the Torus

Consider  $S_{1,0,0}$ : the torus with no punctures or boundaries.

Its mapping class group  $\mathcal{M}_{1,0,0}$  does not involve braids, but we will use it to introduce some new concepts. The great advantage is that  $\mathcal{M}_{1,0,0}$  is easy to understand.

$\mathcal{M}_{1,0,0}$  is generated by Dehn twists (or Lickorish twists).

See Farb+Marg.  
p. 58.



For example, a convenient way of thinking of  $\mathcal{M}_{1,0,0}$  is by looking at its action on  $\pi_1(S_{1,0,0}) = \mathbb{Z}^2$ :

$$D_1, D_2 \in \mathcal{M}_{1,0,0}, \quad \begin{pmatrix} p \\ q \end{pmatrix} \in \pi_1(S_{1,0,0})$$

$$D_1 \begin{pmatrix} p \\ q \end{pmatrix} = \begin{pmatrix} p+q \\ q \end{pmatrix}, \quad D_2 \begin{pmatrix} p \\ q \end{pmatrix} = \begin{pmatrix} p \\ q-p \end{pmatrix}.$$

Here  $p$  and  $q$  just count the # of times the loop winds around each direction of the torus.

So we can write  $D_1$  and  $D_2$  as elements of  $SL_2(\mathbb{Z})$ :

$$D_1 = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}, \quad D_2 = \begin{pmatrix} 1 & 0 \\ -1 & 1 \end{pmatrix},$$

$$D_1^{-1} = \begin{pmatrix} 1 & -1 \\ 0 & 1 \end{pmatrix}, \quad D_2^{-1} = \begin{pmatrix} 1 & 0 \\ 1 & 1 \end{pmatrix}.$$

Note that  
 $D_1^2 D_2 D_1^2 D_2 = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix}$   
 $D_2^2 D_1 D_2^2 D_1 = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix}$

Hence, an arbitrary element of  $M_{1,0,0}$  can be written

$$[h] = \begin{pmatrix} a & b \\ c & d \end{pmatrix}, \quad a, b, c, d \in \mathbb{Z} \\ ad - bc = 1. \quad \leftarrow +1 \text{ to preserve orientation!}$$

Now let us explore the range of behaviour of  $[h]$ .

Consider its characteristic polynomial

$$\det([h] - \lambda I) = \lambda^2 - \underbrace{(\text{trace}[h])}_{a+d} \lambda + \underbrace{\det[h]}_1 \\ = 0 \text{ for an eigenvalue.}$$

Let  $\tau = \text{trace}[h]$ . By the Cayley-Hamilton theorem,

$$[h]^2 - \tau[h] + I = 0, \quad \text{that is, } [h] \text{ satisfies its own char. poly.}$$

The discriminant is  $\tau^2 - 4$ , so we have 3 possible cases:

$$|\tau| < 2, \quad |\tau| = 2, \quad |\tau| > 2.$$

Let us look at each in turn.

1.  $|\tau| < 2$ : Since  $[h] \in SL_2(\mathbb{Z})$ , there are only a few possibilities:  $\tau = 1, \tau = 0, \tau = -1$ .

If  $\tau = 0$ , then  $[h]^2 = -I$ , so  $[h]^4 = I$ .

If  $\tau = \pm 1$ , then  $[h]^2 \mp [h] + I = 0$

$$[h]^2 = \pm [h] - I.$$

$$\begin{aligned} [h]^3 &= \pm [h]^2 - [h] \\ &= \pm (\pm [h] - I) - [h] \\ &= \mp I \end{aligned}$$

So  $[h]^6 = I$ .

So, regardless of the exact value of  $|\tau| < 2$ , we have

$$\boxed{[h]^{12} = I} \quad 12 = \underbrace{2 \cdot 2 \cdot 3}_4^6$$

This shows that in this case  $[h]$  is finite-order: we will define this more precisely later, but for now sufficient to say that  $h \in [h]$  with  $|\tau| < 2$  is such that  $h^m$  is isotopic to id, with  $m > 0$ .

2.  $|\tau| = 2$ :  $\lambda^2 \mp 2\lambda + 1 = 0$   $\lambda = \pm 1$   
with multiplicity 2.  
 $\tau = \pm 2$ .

If  $a = d = \pm 1$ , then  $[h] = \begin{pmatrix} \pm 1 & b \\ c & \pm 1 \end{pmatrix}$ .

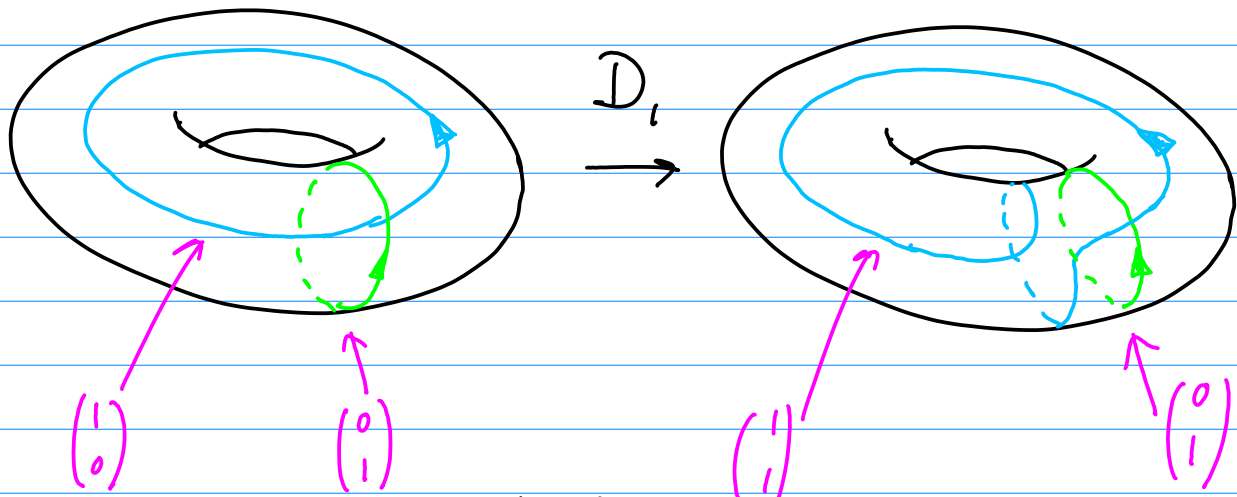
To satisfy  $\det [h] = 1$ , must have  $bc = 0$ , so either  $b$  or  $c$  or both must vanish. If both vanish then  $[h]^2 = I$ , so that  $[h]$  is finite-order again.

Otherwise, take  $b \neq 0$  and  $c = 0$  WLOG.

$[h] = \begin{pmatrix} \pm 1 & b \\ 0 & \pm 1 \end{pmatrix}$ , so  $\begin{pmatrix} 0 \\ 1 \end{pmatrix}$  is an eigenvector with e.v.  $= \pm 1$ .

This eigenvector is in  $\pi_1(T^2)$ . It represents a loop around one of the torus' periodic directions.

Observe in particular that if we take  $\lambda = +1$ , then  $[h] = D_1^b$ , the Dehn twist.



That is:  $D_1$  leaves the loop  $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$  invariant.

Note that this does not imply that all  $h \in [h]$  leaves any curve invariant, but we do know that  $\exists h$  that does.

Now look at the other possibilities:  $a+d = \pm 2$  in general (i.e., w/o  $a=d = \pm 1$ ).

Then note that  $([h] \mp I)^2 = [h]^2 \mp 2[h] + I = 0$   
by Cayley-Hamilton.  
[h]  $\mp$  I is nilpotent, so it is  $\begin{pmatrix} 0 & e \\ 0 & 0 \end{pmatrix}$  in some coordinates.

Hence, there is a change of coordinate  $\underline{x}' = A \underline{x}$   
such that

$$A[h]A^{-1} = \begin{pmatrix} \pm 1 & e \\ 0 & \pm 1 \end{pmatrix} \quad \text{Jordan form}$$

and we are back to the case  $a=d = \pm 1$  above. We conclude once again that a loop (given by the same eigenvector of  $[h]$ ) is invariant.

$$\begin{pmatrix} b \\ \pm 1 - a \end{pmatrix}$$

Conclude: in all cases for  $|\tau| = 2$  the action of  $[h]$  on  $\pi_1(T^2)$  leaves a loop invariant.

(Exception: if  $[h] = \pm I$  then  $[h]$  is finite-order.)

This leaves the most interesting case,  $|\tau| > 2$ .

3.  $|\tau| > 2$ : In that case we get two distinct real roots:

$$\lambda_{\pm} = \frac{1}{2}(\tau \pm \sqrt{\tau^2 - 4}), \text{ with } \lambda_+ \lambda_- = 1$$

The roots are inverse of each other, and we define

$$\lambda = \max(|\lambda_+|, |\lambda_-|) > 1.$$

The eigenvectors of  $[h]$  are  $\pm = \text{sign}(\tau)$

$$u = \begin{pmatrix} \pm\lambda - d \\ c \end{pmatrix}, \quad s = \begin{pmatrix} \pm\lambda^{-1} - d \\ c \end{pmatrix}, \quad \lambda > 1$$

Check: (with  $\tau > 2$ )

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} \lambda - d \\ c \end{pmatrix} = \begin{pmatrix} a\lambda - ad + bc \\ c\lambda - cd + cd \end{pmatrix} = \begin{pmatrix} (-\lambda^2 + (a+d)\lambda - 1) + \lambda^2 - \lambda d \\ c\lambda \end{pmatrix} = \lambda \begin{pmatrix} \lambda - d \\ c \end{pmatrix}.$$

Claim:  $\lambda$  is irrational.

Assume  $\tau > 2$  (positive): ( $\leftarrow \tau < -2$  is of course almost the same proof)

$$\lambda = \frac{1}{2}(\tau + \sqrt{\tau^2 - 4}) = a_0 + \frac{1}{a_1 + \frac{1}{a_2 + \frac{1}{\dots}}}$$

$$\lambda^2 - \tau\lambda + 1 = 0 \iff \lambda - \tau + \lambda^{-1} = 0$$

$$\lambda = \tau - 1 + (1 - \lambda^{-1}) = a_0 + (1 - \lambda^{-1})$$

Positive integer

$$0 < 1 - \lambda^{-1} < 1$$

Fractional part

$$\therefore a_0 = \tau - 1$$

$$\lambda - a_0 = 1 - \lambda^{-1} = \frac{1}{a_1 + \frac{1}{a_2 + \frac{1}{\dots}}}$$

$$\frac{1}{1 - \lambda^{-1}} = a_1 + \frac{1}{a_2 + \frac{1}{\dots}}$$

$$1 - \lambda^{-1} = 1 - (\tau - \lambda) \quad \text{since } \lambda^2 - \tau\lambda + 1 = 0$$

$$= (1 - \tau) + \lambda$$

$\phi^2!$

Since  $\lambda = \frac{1}{2}(\tau + \sqrt{\tau^2 - 4})$  is monotonic,  $\min \lambda = \frac{1}{2}(3 + \sqrt{5})$   
for  $|\tau| = 3$ .

Hence,  $1 < \frac{1}{1 - \lambda^{-1}} \leq \frac{1}{1 - \max \lambda^{-1}} = 2.61803\dots$

$$= \frac{1}{\frac{1}{2}(\sqrt{5} - 1)}$$

$$= \frac{2(\sqrt{5} + 1)}{4} = \frac{1}{2}(1 + \sqrt{5}) < 2.$$

$\therefore 1 < \frac{1}{1 - \lambda^{-1}} < 2$ . Conclude:  $a_1 = 1$ .

$$\left(\frac{1}{1 - \lambda^{-1}} - 1\right)^{-1} = a_2 + \frac{1}{a_3 + \frac{1}{\dots}}$$

$$\left(\frac{1}{1 - \lambda^{-1}} - 1\right)^{-1} = \left(\frac{\lambda - (\lambda - \lambda^{-1})}{1 - \lambda^{-1}}\right)^{-1} = \left(\frac{1}{\lambda - 1}\right)^{-1} = \lambda - 1.$$

$$\lambda - 1 = a_2 + \frac{1}{a_3 + \frac{1}{\dots}} \iff \lambda = \underbrace{(a_2 + 1)}_{a_0 = \tau - 1} + \frac{1}{a_3 + \frac{1}{\dots}}$$

Same expression as when we started!

$$a_0 = \tau - 1$$

$$\therefore a_2 = \tau - 2.$$

One more time:

$$(\lambda - 1 - a_2)^{-1} = a_3 + \frac{1}{a_4 + \frac{1}{\dots}}$$

$$\lambda^2 - \tau\lambda + 1 = 0$$

$$\lambda - \tau + \lambda^{-1} = 0$$

$$\lambda - \tau + 2 = 2 - \lambda^{-1}$$

$$(\lambda - 1 - a_2)^{-1} = (\lambda - \tau + 1)^{-1} = (1 - \lambda^{-1})^{-1}$$

$$\frac{1}{1 - \lambda^{-1}} = a_3 + \frac{1}{a_4 + \frac{1}{\dots}} \quad \therefore \begin{aligned} a_3 &= a_1 \\ a_4 &= a_2 \\ a_5 &= a_1 \\ a_6 &= a_2 \dots \end{aligned}$$

→ We've seen this before!

... Periodic!

Continued fraction representation:

$$\lambda = [\tau - 1; \underbrace{1, \tau - 2, 1, \tau - 2, \dots}_{\text{period-2}}]$$

$\therefore$  So  $\lambda$  is irrational.

Of course, this shouldn't have surprised us: the irrational solutions of quadratic equations with integer coefficients are periodic. But this polynomial has a particularly simple form.

More importantly, this shows that  $\lambda$  is irrational for any  $|\tau| > 2$  (negative  $\tau$  proceeds identically)

$[h]$  is the isotopy class of an Anosov diffeomorphism with dilatation  $\lambda > 1$ .

More precisely, in the isotopy class  $[h]$  there is a unique diffeo  $h_A : T^2 \rightarrow T^2, T^2 = [X, Y]$ , defined by

$$h_A(\underline{r}) = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} \text{ mod } 1 \quad \underline{r} = \begin{pmatrix} x \\ y \end{pmatrix} \in T^2 = [X, Y]$$

$|a+d| > 2.$

$h_A$  is an Anosov diffeomorphism.

A leaf of the unstable foliation  $\mathcal{F}^u$  of  $h_A$  is generated by

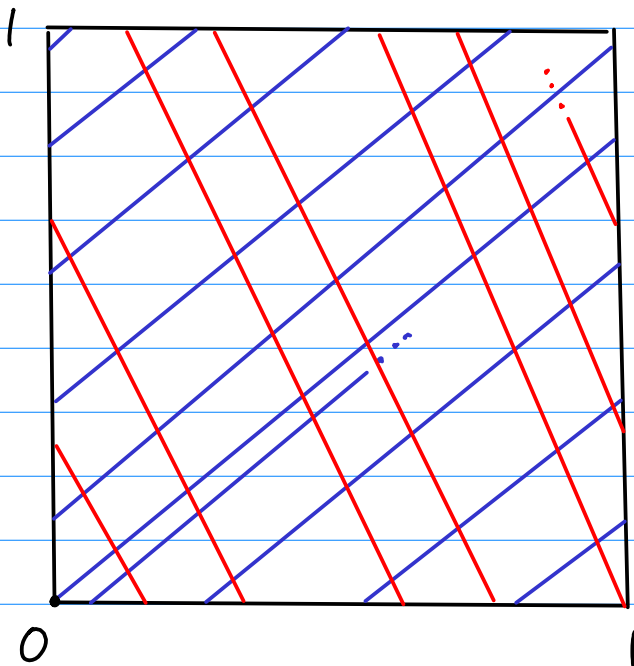
$$h_A^m(\pm u) = \lambda^m (\pm u) \text{ mod } 1, \quad m \in \mathbb{Z}, m > 0, \quad u = \begin{pmatrix} \pm\lambda - d \\ c \end{pmatrix}$$

$\downarrow$  regarded as a curve.

Similarly, a leaf of the stable foliation  $\mathcal{F}^s$  of  $h_A$  is generated by

$$h_A^{-m}(\pm s) = \lambda^{-m} (\pm s) \text{ mod } 1, \quad m \in \mathbb{Z}, m > 0, \quad s = \begin{pmatrix} \pm\lambda - d \\ c \end{pmatrix}$$

The leaves of the foliations do not carry an orientation.



Both leaves wind around the torus  $\infty$  many times, since their slope is irrational. Each leaf is dense in  $T^2$ .

They are transverse, that is they are nowhere tangent.

The foliation  $\mathcal{F}^{u,s}$  can be generated by considering the iterates of all possible curves with slope parallel to  $u, s$ .

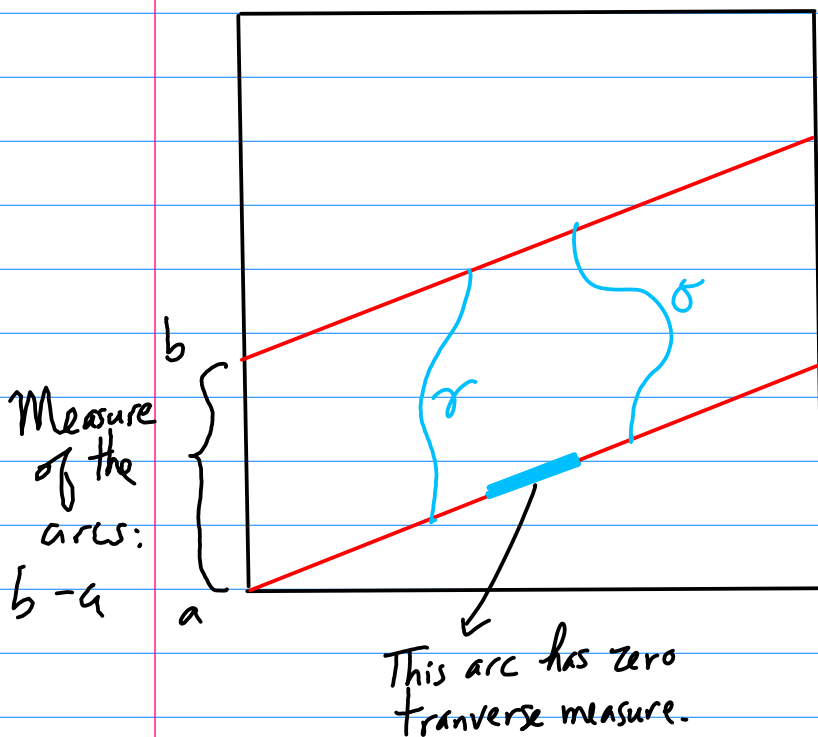
Observe that the foliations are invariant under  $h_A$ , since they have the same slope as its eigenvectors.

If we take an arbitrary curve in  $T^2$ , then clearly it will converge to  $\mathcal{F}^u$  under repeated application of  $h_A$ , or to  $\mathcal{F}^s$  under repeated application of  $h_A^{-1}$ .

We can turn our foliations into measured foliations by equipping them with transverse measures  $\mu^u$  and  $\mu^s$ .

Obtain measured foliations  $(\mathcal{F}^u, \mu^u)$  and  $(\mathcal{F}^s, \mu^s)$ .

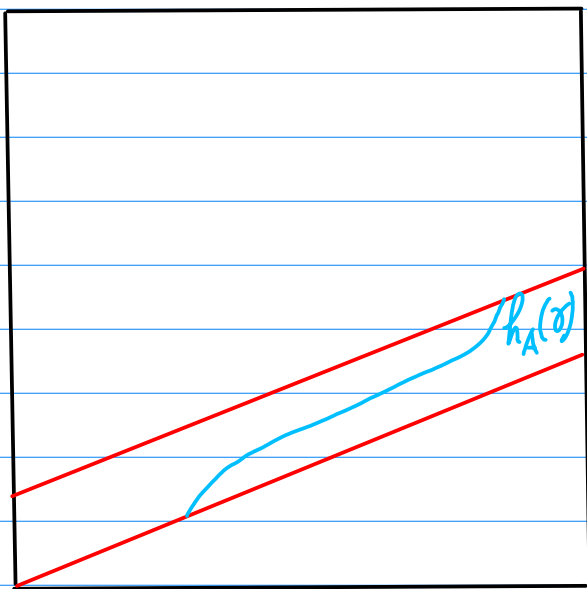
The measures are the usual Borel measures, taken as proportional to the distance between leaves.



We define the transverse measure of arcs transverse to  $\mathcal{F}^u$  (not tangent to any leaf) by, for instance, the standard Borel measure as on the left. The arcs  $\tau$  and  $\sigma$  have the same measure.

$$\mu^u(\tau) = \mu^u(\sigma).$$

The important thing is what happens to the measure under the action of  $h_A$ .



The arc  $h_A(r)$  now has endpoints on different leaves, though it remains transverse. In fact,

$$\mu^u(h_A(r)) = \lambda^{-1} \mu^u(r)$$

If we define the action of  $h_A$  on  $\mu^u$  by.

$$h_{A*}(\mu^u)(r) = \mu^u(h_A(r)) = \lambda^{-1} \mu^u(r),$$

or

$$h_{A*}(\mu^u) = \lambda^{-1} \mu^u.$$

Similarly, we have  $h_{A*}(\mu^s) = \lambda \mu^s$ .

Hence, altogether we write

$$h_{A*}(\mathcal{F}^u, \mu^u) = (\mathcal{F}^u, \lambda^{-1} \mu^u)$$

$$h_{A*}(\mathcal{F}^s, \mu^s) = (\mathcal{F}^s, \lambda \mu^s)$$

← "Dilatation"

This last claim requires a bit more explanation:

$$h_A(\mathcal{F}^u, \mu^u) = (\mathcal{F}^u, \lambda \mu^u) \quad \text{and expr. for stable foliation.}$$

↑  
why not  $\lambda^{-1}$ ?

The new foliation is just  $h_A(\mathcal{F}^u) = \mathcal{F}^u$  (invariant)

We define the measure on  $h_A(\mathcal{F}^u)$  by pulling back arcs  $\gamma$  transverse to  $h_A(\mathcal{F}^u)$ : but in general it might not be!

$$h_A^*(\mu^u)(\gamma) = \mu^u(h_A^{-1}(\gamma)) = \lambda \mu^u(\gamma)$$

The measure of a transverse arc increases under  $h_A^{-1}$ .

So  $h_A(\mathcal{F}^u, \mu^u) = (\mathcal{F}^u, \lambda \mu^u)$

Similarly,  $h_A(\mathcal{F}^s, \mu^s) = (\mathcal{F}^s, \lambda^{-1} \mu^s)$ ,

since  $h_A^*(\mu^s)(\gamma) = \mu^s(h_A^{-1}(\gamma)) = \lambda^{-1} \mu^s(\gamma)$ .

Another way to look at measured foliations:

Define the foliation  $\mathcal{F}^u$  in terms of the closed 1-form

$$\omega^u = dy - m^u dx$$

where  $m$  is the slope of the unstable direction:

$$m^u = \frac{c}{\lambda - d} \quad (\tau = a + d > 2) \quad \left( m^s = \frac{c}{\lambda^{-1} - d} \right)$$

Then we define the transverse measure  $\mu^u(\gamma)$  of an arc  $\gamma$  by

$$\mu^u(\gamma) = \int_{\gamma} \omega^u = \int_{\gamma} dy - m^u dx$$

We can then think of leaves of  $\mathcal{F}^u$  as consisting of lines of vanishing  $\omega^u$ :

$$d(\underbrace{m^u x + b}_{y(x)}) - m^u dx = 0.$$

The action of  $h_A$  on  $\mu^u$  can then be calculated directly:

$$\mu^u(h_A(\gamma)) = \int_{h_A(\gamma)} \omega^u(\underline{r}) = \int_{h_A(\gamma)} dy - m^u dx$$

Claim:  $\omega^u(h_A(r)) = \lambda^{-1} \omega^u(r)$   $\lambda^2 - (a+d)\lambda + 1 = 0$

Recall  $h_A(r) = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$ ,  $\lambda = (a+d) - \lambda^{-1}$

so  $\omega^u(h_A(r)) = d(cx + yd) - m^u d(ax + by)$   
 $= (c - m^u a) dx + (d - m^u b) dy$

$c - m^u a = c \left( 1 - \frac{a}{\lambda - d} \right) = \left( \frac{\lambda - (a+d)}{\lambda - d} \right) c = -\lambda^{-1} m^u$

$d - m^u b = d - \frac{bc}{a - \lambda^{-1}} = \frac{d\lambda - (ad - bc)}{a - \lambda^{-1}} = \lambda^{-1} \frac{d - \lambda}{a - \lambda^{-1}} = \lambda^{-1}$

Hence,  $\omega^u(h_A(r)) = \lambda^{-1} (dy - m^u dx)$   
 $= \lambda^{-1} \omega^u(r)$

This means that  $\mu^u(h_A(r)) = \int_{h_A(r)} \omega^u(r) = \int_r \omega^u(h_A(r))$   
 $= \lambda^{-1} \mu^u(r),$

as before.