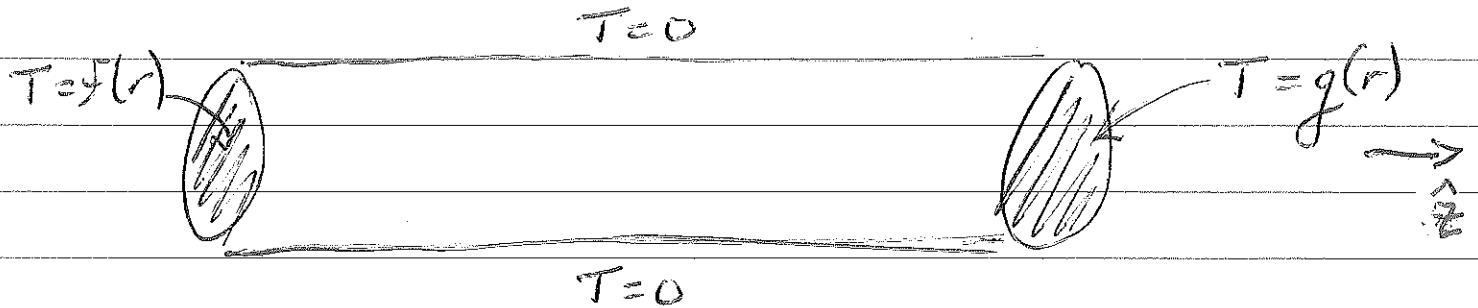


Another physical context (simpler, Bessel of order zero)

$\frac{\partial T}{\partial t} = k \nabla^2 T$ in cylindrical coordinates and the steady-state limit



$T(\theta, z)$ bounded $\Rightarrow f(r), g(r)$ bounded as $r \rightarrow 0$
and for consistency we take
 $f(r) = g(r) = 0$ at $r=1$

In cylindrical coordinates ; steady state

$$\frac{\partial T}{\partial t} = k \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial z^2} \right) T$$

If we also consider axisymmetric solutions :

$$\left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{\partial^2}{\partial z^2} \right) T = 0$$

New separation of variables $T(r, z) = \varphi(r) \psi(z)$

$$\Rightarrow \frac{1}{\varphi} \left[\varphi'' + \frac{1}{r} \varphi' \right] = -\frac{\psi''}{\psi} = -\lambda$$

$$\Rightarrow \varphi''(r) + \frac{1}{r} \varphi'(r) + \lambda \varphi(r) = 0 \quad \text{Bessel of order zero}$$

interior $\varphi(0)$ banded

sidewalls $\varphi(1) = 0$

$$\psi''(z) - \lambda \psi(z) = 0$$

endwalls $T(r, 0) = f(r) = \varphi(r) \psi(0)$

$T(r, 1) = g(r) = \varphi(r) \psi(1)$

Bessel of order zero in SL form

$$\frac{d}{dr} \left(r \frac{d\varphi}{dr} \right) = -\lambda r \varphi(r)$$

$\varphi(0)$ banded, $\varphi(1) = 0$

How do we solve Bessel's Eqn? (a 2nd-order linear, non-constant coefficient, homogeneous ODE)

For convenience, let's scale out λ :

$$\text{let } x = \sqrt{\lambda} r, \quad \frac{dx}{dr} = \sqrt{\lambda}$$

$$\frac{dY}{dr} = \frac{dY}{dx} \frac{dx}{dr} = \sqrt{\lambda} \frac{dY}{dx}$$

$$\frac{d^2Y}{dr^2} = \frac{d}{dr} \left(\sqrt{\lambda} \frac{dY}{dx} \right) = \sqrt{\lambda} \frac{d^2Y}{dx^2} \frac{dx}{dr} = \lambda \frac{d^2Y}{dx^2}$$

$$\lambda \frac{d^2Y}{dx^2} + \frac{\sqrt{\lambda}}{x} \sqrt{\lambda} \frac{dY}{dx} + \lambda Y = 0 \quad Y = Y(x)$$

$$\Rightarrow \left[\begin{array}{l} y'' + \frac{1}{x} y' + y = 0 \\ y(0) \text{ bounded} \\ y(\sqrt{\lambda}) = 0 \end{array} \right]$$

This type of equation has a Frobenius series soln.

$$y(x) = x^\alpha \sum_{n=0}^{\infty} C_n x^n \quad (\text{about } x_0 = 0)$$

So-called "standard form" (not "SL" Form): ⑥

$$y'' + \tilde{p}(x)y' + \tilde{q}(x)y = 0$$

Homogeneous, linear, non-constant coefficient

* For $\tilde{p}(x), \tilde{q}(x)$ analytic in a neighborhood of x_0 in the complex plane (x_0 is an ordinary point)

$\Rightarrow \tilde{p}(x), \tilde{q}(x)$ have Taylor Series expansions

$$\tilde{p}(x) = \sum_{n=0}^{\infty} p_n (x-x_0)^n \quad n \text{ integer}$$

$$\tilde{q}(x) = \sum_{n=0}^{\infty} q_n (x-x_0)^n \quad n \text{ integer}$$

Then the solution has a convergent Taylor series expansion

$$y(x) = \sum_{n=0}^{\infty} a_n (x-x_0)^n \quad n \text{ integer}$$

with radius of convergence at least as large as the distance from x_0 to the nearest singularity in the complex plane

* For x_0 a regular singular point :

$$(x-x_0)^2 \tilde{p}(x) = \sum_{n=0}^{\infty} p_n (x-x_0)^n \quad n \text{ integer}$$

$$(x-x_0)^2 \tilde{q}(x) = \sum_{n=0}^{\infty} q_n (x-x_0)^n \quad n \text{ integer}$$

Then there is at least one solution of the form

$$y(x) = (x-x_0)^\alpha \sum_{n=0}^{\infty} a_n (x-x_0)^n \quad \begin{array}{l} \alpha \text{ real} \\ n \text{ integer} \end{array}$$

$$= (x-x_0)^\alpha A(x)$$

where $A(x)$ is analytic at $x=x_0$, with Taylor series expansion at least as large as the distance to the nearest singularity in the complex plane.

Back to Bessel of order zero:

"standard form" $y'' + \frac{1}{x} y' + y = 0$

$y(0), y'(0)$ bounded $y(\sqrt{1}) = 0$

$x_0 = 0$ is a regular singular point:

$x^2 p(x) = x^0 \frac{1}{x} = \frac{1}{x}$ analytic

$x^2 q(x) = x^2$ analytic

Therefore we look for a solution

$y(x) = x^\alpha \sum_{n=0}^{\infty} a_n x^n = \sum_{n=0}^{\infty} a_n x^{n+\alpha}$

$y'(x) = \sum_{n=0}^{\infty} (\alpha+n) a_n x^{n+\alpha-1}$

$y''(x) = \sum_{n=0}^{\infty} (\alpha+n)(\alpha+n-1) a_n x^{n+\alpha-2}$

and plug into $x^2 y'' + x y' + x^2 y = 0$

(easier if we don't need to divide)

$$\sum_{n=0}^{\infty} a_n (\alpha+n)(\alpha+n-1) x^{\alpha+n} + \sum_{n=0}^{\infty} a_n (\alpha+n) x^{\alpha+n} + \sum_{n=0}^{\infty} a_n x^{\alpha+n+2} = 0$$

Equate like powers of x :

$$x^{\alpha} : a_0 \alpha (\alpha-1) + a_0 \alpha = 0$$

For $a_0 \neq 0 \Rightarrow \alpha^2 = 0 \Rightarrow \alpha = 0$ repeated root

$$x^{\alpha+1} : a_1 (\alpha+1) \alpha + a_1 (\alpha+1) = 0$$

$$a_1 (\alpha^2 + 2\alpha + 1) = 0 \Rightarrow a_1 = 0$$

$$x^{\alpha+2} : a_2 (\alpha+2)(\alpha+1) + a_2 (\alpha+2) + a_0 = 0$$

$$x^{\alpha+3} : a_3 (\alpha+3)(\alpha+2) + a_3 (\alpha+3) + a_1 = 0$$

etc. \Rightarrow all odds zero and

$$a_n(\alpha) = - \frac{a_{n-2}}{(\alpha+n)^2} \quad n \geq 2$$

with $\alpha=0 \Rightarrow$

$$y_1(x) = a_0 \left[1 + \sum_{m=1}^{\infty} \frac{(-1)^m x^{2m}}{2^{2m} (m!)^2} \right] = a_0 J_0(x)$$

The Bessel Function of the 1st kind of order zero

Since our eqn. is 2nd-order, we need to find a 2nd linearly independent solution

* use Reduction of Order, or

* By analogy with Euler equations, try

$$y_2(x) = J_0(x) \ln x + x^\alpha \sum_{n=0}^{\infty} b_n x^n$$

↑ different coefficients

$$y_2(x) = J_0(x) \ln x + \sum_{m=1}^{\infty} \frac{(-1)^{m+1} H_m x^{2m}}{2^{2m} (m!)^2}$$

$$H_m = 1 + \frac{1}{2} + \frac{1}{3} + \dots + \frac{1}{m}$$

Back to our pipe flow problem, the solution for the radial dependence is

$$\psi(r) = C_1 J_0(\sqrt{\lambda} r) + C_2 y_2(\sqrt{\lambda} r)$$

since $x = \sqrt{\lambda} r$

But $y_2(\sqrt{\lambda} r)$ has a factor of $\ln(\sqrt{\lambda} r)$ so it is not bounded at $r=0$
 $\Rightarrow C_2 = 0$

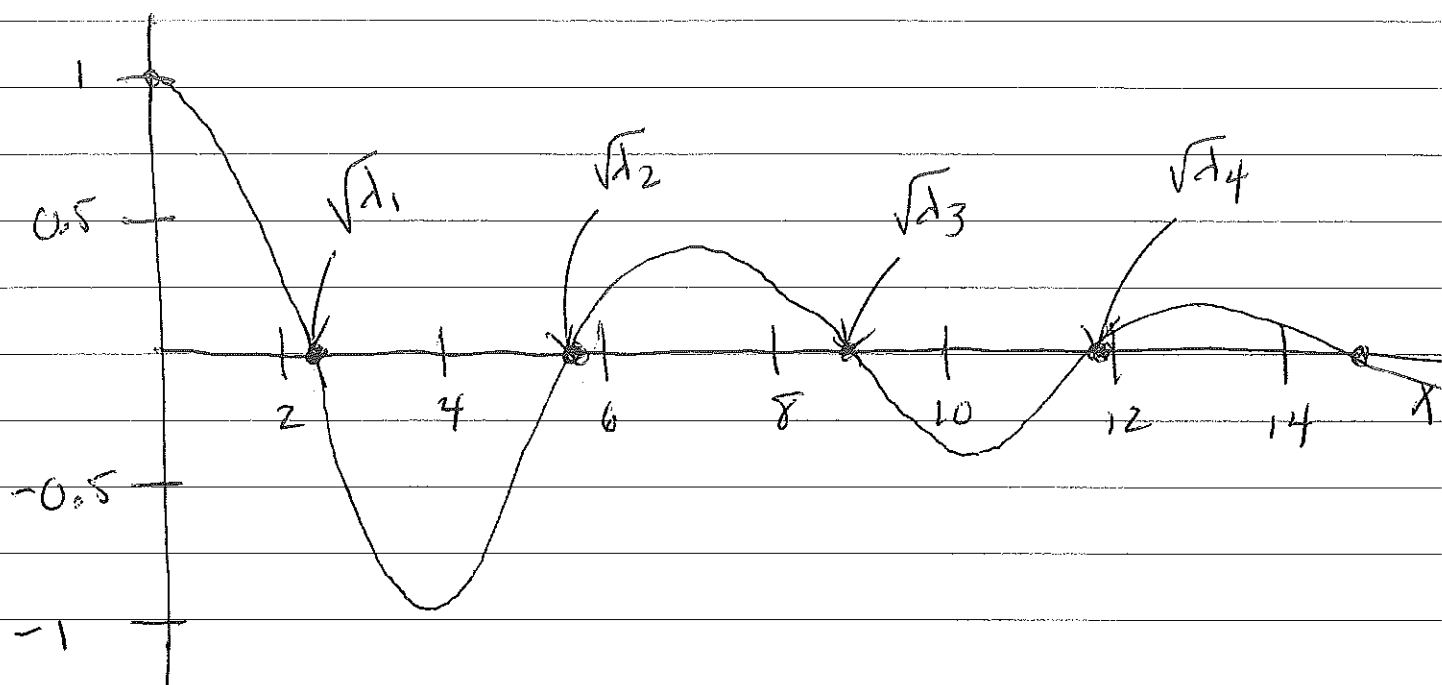
$$\Rightarrow \psi(r) = C_1 J_0(\sqrt{\lambda} r) \text{ with } \psi(1) = 0$$

and the boundary condition $\psi(1) = 0$

determines the values of λ :

$$\psi(1) = J_0(\sqrt{\lambda}) = 0$$

(9)

 $J_0(x)$ 

$\lambda_1 < \lambda_2 < \lambda_3 < \dots < \lambda_n < \dots$
the eigenvalues

$$\lim_{n \rightarrow \infty} \lambda_n = \infty$$

$$J_0(x) \approx \left(\frac{2}{\pi x} \right)^{1/2} \cos\left(x - \frac{\pi}{4}\right) \text{ as } x \rightarrow \infty$$

For each λ_n , there is an eigenfunction

$$\psi_n(r) = J_0(\sqrt{\lambda_n} r) \text{ that satisfies the}$$

ODE and the boundedness condition at $r=0$

and $\psi_n(1) = 0$