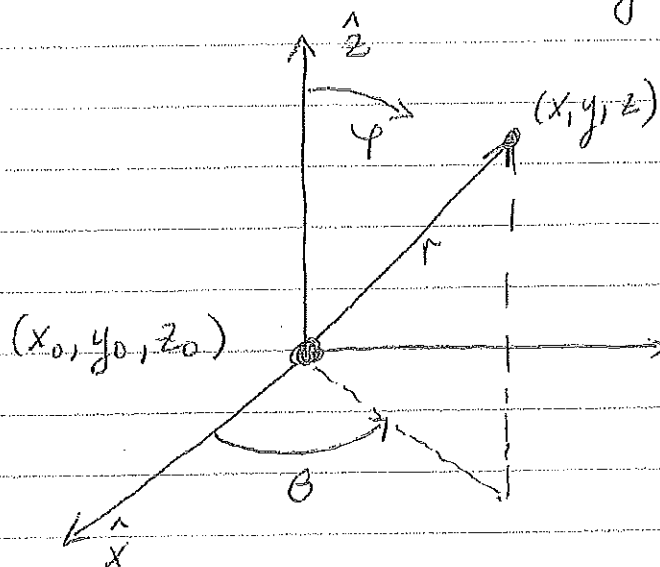


# Math 322 | The infinite space Green's function

for use in the Method of Images:



We need to solve  $\nabla^2 G(\underline{x}; \underline{x}_0) = \delta(\underline{x} - \underline{x}_0)$

and the idea is that the effect of this "source" at other values of  $\underline{x}$  should depend only

on  $r$ , but not on  $\theta, \varphi$ , where

$$r = \left[ (x-x_0)^2 + (y-y_0)^2 + (z-z_0)^2 \right]^{1/2}$$

$$\Rightarrow \nabla^2 G(\underline{x}; \underline{x}_0) = \nabla^2 G(r)$$

with  $\nabla^2$  in spherical coordinates given by

(2)

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{1}{\sin^2 \varphi} \frac{\partial^2}{\partial \theta^2} + \frac{1}{r^2} \frac{1}{\sin \varphi} \frac{\partial}{\partial \varphi} \left( \sin \varphi \frac{\partial}{\partial \varphi} \right)$$

$$\Rightarrow \frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{dG(r)}{dr} \right) = 0 \quad \underline{x} \neq \underline{x}_0$$

$$\Rightarrow r^2 \frac{dG}{dr} = C \quad \Rightarrow G(r) = -\frac{C}{r} + D$$

Now lets integrate around the "source" :

$$\int_{\mathcal{V}} \nabla^2 G d\mathcal{V} = \int_{\mathcal{V}} \delta(\underline{x} - \underline{x}_0) d\mathcal{V} = 1$$

$$\int_{\mathcal{V}} \nabla \cdot \nabla G d\mathcal{V} = \int_A \nabla G \cdot \hat{n} dA$$

and lets consider a sphere at fixed  $r$  :

$\nabla G$  is in the  $\hat{r}$ -direction

$\hat{n}$  is in the  $\hat{r}$ -direction

(3)

$$dA = r^2 \sin \varphi d\varphi d\theta \quad \text{at fixed } r$$

$$\int_0^{2\pi} d\theta \int_0^{\pi} \sin \varphi d\varphi \quad r^2 \frac{dG(r)}{dr} = 1$$

$$= r^2 \frac{dG}{dr} \left[ -\cos \varphi \Big|_0^{\pi} \right] \theta \Big|_0^{2\pi} = 1$$

$$= 2 \cdot 2\pi r^2 \frac{dG}{dr} = 1 \quad \text{or} \quad \frac{dG}{dr} = \frac{1}{4\pi r^2}$$

and the constant  $C = \frac{1}{4\pi} \Rightarrow$

$G(r) = -\frac{1}{4\pi r} + D$	determined up to an arbitrary constant $D$
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Green's Function for other linear operators.

Consider the 1D Heat Eqn.

$$\frac{du}{dt} = K \frac{d^2u}{dx^2} + \phi(x, t) \quad \begin{array}{l} 0 < x < L \\ t > 0 \end{array}$$

$$u(x, 0) = F(x)$$

$u(0, t)$ ,  $u(L, t)$  given

or  $\frac{du}{dx}(0, t)$ ,  $\frac{du}{dx}(L, t)$  given

or mixed boundary conditions

Multiply by a Green's Function and integrate "over the whole domain" which now includes time

$$\int_0^L dx \int_0^{\infty} dt G(x, t; x_0, t_0) \left[ \frac{du}{dt}(x, t) \right]$$

$$= K \frac{d^2u(x, t)}{dx^2} + \phi(x, t) \Big]$$

Now integrate by parts:

$$(i) \int_0^{\infty} G \frac{du}{dt} dt = - \int_0^{\infty} u \frac{dG}{dt} dt + uG \Big|_{t=0}^{\infty}$$

$$(ii) \int_0^L G \frac{d^2 u}{dx^2} dx = \int_0^L u \frac{d^2 G}{dx^2} dx - \left[ u \frac{dG}{dx} - G \frac{du}{dx} \right] \Big|_{x=0}^L$$

So we arrive at

$$\begin{aligned} & \int_0^L \int_0^{\infty} \left( -u \frac{dG}{dt} - k u \frac{d^2 G}{dx^2} \right) dx dt \\ &= \int_0^L \int_0^{\infty} G G dx dt - \int_0^L dx uG \Big|_{t=0}^{\infty} \\ & \quad - k \int_0^{\infty} dt \left[ u \frac{dG}{dx} - G \frac{du}{dx} \right] \Big|_{x=0}^L \end{aligned}$$

(3)

Now choose  $-\frac{\partial G}{\partial t} - k \frac{\partial^2 G}{\partial x^2} = \delta(x-x_0)\delta(t-t_0)$

$$G = G(x, t; x_0, t_0)$$

so that the LHS =  $u(x_0, t_0)$

where  $0 < x_0 < L$ ,  $0 < t_0 < \infty$

Note that  $G(x, t; x_0, t_0)$  satisfies a backwards heat equation

$$-\frac{\partial G}{\partial t} = k \frac{\partial^2 G}{\partial x^2} + \delta(x-x_0)\delta(t-t_0)$$

so time is in the opposite sense and the analogy to the homogeneous condition in time is

$$G(x, t; x_0, t_0) = 0 \text{ for } t > t_0$$

Choose boundary conditions as usual

Then what happens to the formula:

$$u(x_0, t_0) = \int_0^L dx \int_0^{t_0} G(x, t) Q(x, t) \quad \leftarrow \begin{array}{l} \text{non-homogeneous} \\ \text{term in} \\ \text{the eqn.} \end{array}$$

$$+ \int_0^L dx F(x) G(x, t) \Big|_{t=0} \quad \leftarrow \begin{array}{l} \text{effects of the} \\ \text{initial condition} \end{array}$$

$$- K \int_0^{\infty} dt \left[ u \frac{dG}{dx} - G \frac{du}{dx} \right] \Big|_{x=0}^{x=L} \quad \leftarrow \begin{array}{l} \text{boundary} \\ \text{conditions} \end{array}$$

So we need to solve

$$-\frac{\partial G}{\partial t} = k \frac{\partial^2 G}{\partial x^2} + \delta(x-x_0)\delta(t-t_0) \quad 0 < t < t_0$$

$$0 < x < L$$

\* homogeneous boundary conditions in space

\*  $G=0 \quad t > t_0$

If we want to use the method of images with the infinite space Green's function

$$-\frac{\partial G}{\partial t} = k \frac{\partial^2 G}{\partial x^2} + \delta(x-x_0)\delta(t-t_0) \quad 0 < t < t_0$$

$$-\infty < x < \infty$$

\*  $G(\pm \infty, t) = 0$

\*  $G(x, t) = 0 \quad t > t_0$

$$G(x, t) = \frac{1}{\sqrt{4\pi k(t_0-t)}} \exp\left[-\frac{(x-x_0)^2}{4k(t_0-t)}\right]$$

$0 < t < t_0$



Example

$$u_t = K u_{xx} + Q, \quad 0 \leq x \leq L, \quad 0 < t < \infty$$

$$u(0,t) = u(L,t) = 0 \quad u(x,0) = F(x)$$

Use the method of images with infinite-space Green's function:

The boundary conditions  $G(0,t) = G(L,t) = 0$  will be satisfied if we take

\* positive sources at  $x = x_0 + 2Ln$

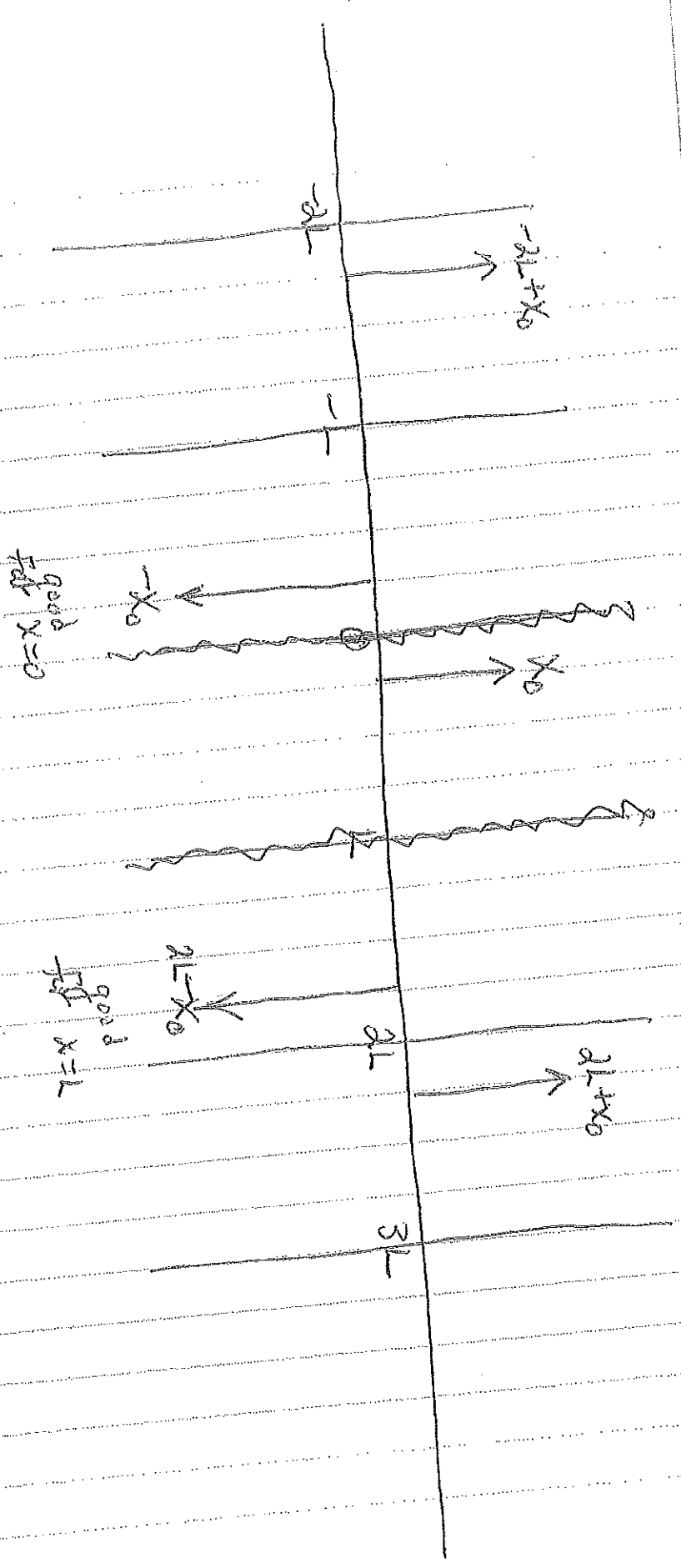
$$n \text{ integer } -\infty < n < \infty$$

\* negative sources at  $x = -x_0 + 2Ln$

$$n \text{ integer } -\infty < n < \infty$$

Need positive scores at  $x_0, x_0 + \Delta L, x_0 + 2\Delta L, -2\Delta L - x_0, -4\Delta L - x_0, \text{etc.}$

Need negative scores at  $-x_0, -\Delta L - x_0, \Delta L - x_0, 3\Delta L - x_0, \text{etc.}$



$$G(x_0, t_0; x, t) = \frac{1}{\sqrt{4\pi K(t-t_0)}} \sum_{n=-\infty}^{\infty} \left. \right\}$$

$$\exp \left[ -\frac{(x-x_0-dnL)^2}{4K(t-t_0)} \right]$$

$$- \exp \left[ -\frac{(x+x_0-dnL)^2}{4K(t-t_0)} \right] \left. \right\} \quad t > t_0$$

$$G(x_0, t_0; x, t) = 0 \quad t < t_0$$